



# EFFECT OF "THERMAL EXPLOSION"

**Lidia Egorova**





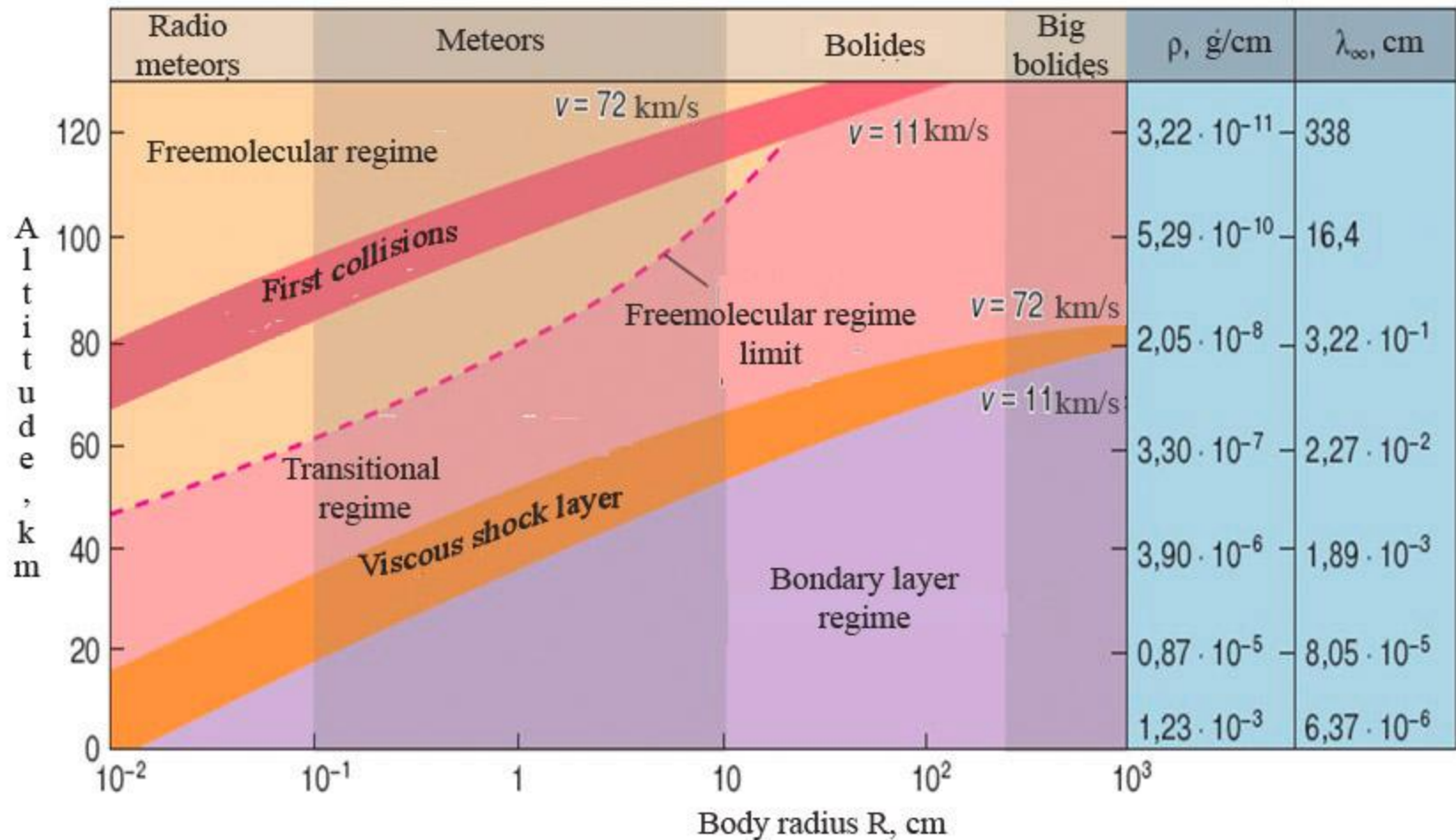
# Spacecraft and meteoroid

- The shape is known and unchanged
- The material is known
- The mass is known and constant
- $V < 16$  km/s (usually 5 - 7 km/s)
- The shape is unknown and can vary
- The material is unknown
- The mass is unknown and can vary
- $V < 72$  km/s (usually 11 - 72 km/s)



# Meteoroid undergoes

- Aerodynamic drag
- Mechanics loading
- Aerodynamic heating
- Radiation
- Ablation
- Fragmentation
- Explosion



*Tirskii, G.A.*, Interaction between Cosmic Bodies and the Earth's and Planet's Atmosphere, Soros Educational Journal, v.6, 2, 2000 (in Russian)



# Physical theory of meteors

$$M\dot{V} = -\frac{1}{2}AC_D\rho V^2, \quad Q\dot{M} = -\frac{1}{2}AC_H\rho V^3,$$

***M*** – mass,

***V*** – velocity,

***A*** – square middle section ,

**$\rho$**  – air density,

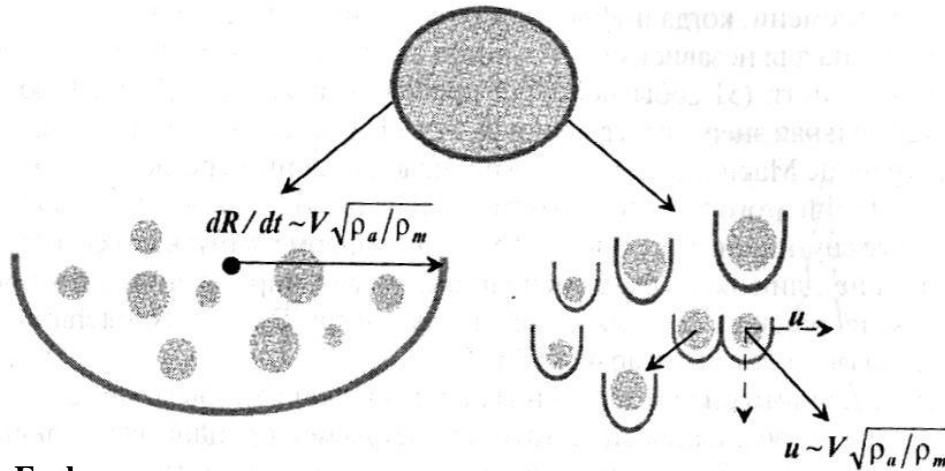
***CH*** – heat transfer coefficient,

***Q*** – specific heat of ablation

(effective enthalpy)



# Modern models of fracture meteoroids



## 1. «Early»

- Jaccia, 1955; Levin, 1962-63; Simonenko 1973-74 – fragments separation from the main body.
- Stanukovich and Shalimov, 1961; ReVelle, 1979 – because of the aerodynamic heating.
- Cook, 1960; Lebedinets and Portnyagin, 1967 – unstable liquid drop.
- McCrosky and Ceplecha, 1969 – because of the previous cosmic collisions.
- Petrov and Stulov, 1976 – because of extremely low density of the body.

## 2. Catastrophic

- Pokrovsky, 1964-66 – “explosion”.
- Jones and Kaiser, 1967; McCrosky and Ceplecha, 1969 – thermal shock.
- Fujiwara, 1989; Jenniskens, 1994 – instantaneous disintegration into fragments.
- ReVelle, 1999-2001 – modification of a single body model.

## 3. Hydrodynamic

- Grigorian, 1976-79; Bronshten, 1985 – the propagation of fragmentation front.
- Hills and Goda, 1993 – the expansion of debris cloud.
- Korobeinikov, Chushkin, Shurshalov, 1986-98 – the fragmentation from inside.
- Melosh, 1981; V.A.Ivanov, 1986-88 – the exploitation of body geometry specifics.
- Chyba, Thomas, Zahnle, 1992-93 – the fluttering of cylindrical body.
- Svetsov, Nemchinov, Teterev, 1995; Shuvalov, Artem’eva, Trubetskaya, 2000 – numerical modeling.

## 4. Progressive, discrete

- Fadeenko, 1967 – comparison of body strength with aerodynamic load.
- Baldwin and Sheaffer, 1971; Tsvetkov, Skripnik 1991; Stulov, 1998; Artem’eva, Shuvalov, 2001 – the application of statistical theory of strength (Weibull, 1939).
- Ivanov, Ryzhansky, 1995 – successive doubling of fragments.
- Ceplecha, Spurny, Borovička, and Keclikova, 1993 – gross-fragmentation.



# **Thermal explosion**

**is a very fast energy transfer to the atmosphere with the conversion of the whole body mass into vapor**



# Elastic theory equation set

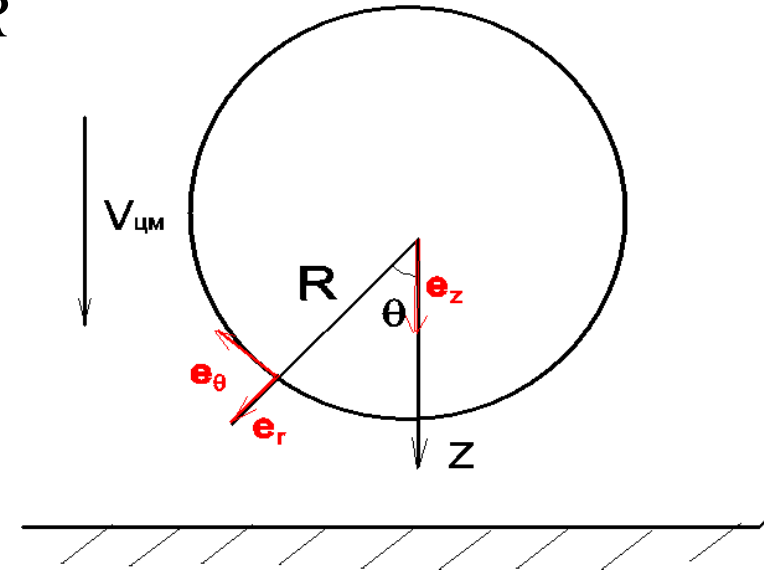
$$\left\{ \begin{array}{l} \frac{1}{1-2\nu} \text{grad div } \mathbf{u} + \Delta \mathbf{u} = -\frac{1}{8G} (\alpha + \beta) \rho V^2 \frac{1}{R} \mathbf{k} \\ \sigma_r = \sigma(\theta), \quad \text{if } r = R \\ \tau_{r\theta} = \tau(\theta), \quad \text{if } r = R \end{array} \right.$$

$\sigma_r$  and  $\tau_{r\theta}$  are from aerodynamic solution of hypersonic flow

$\nu$  - Poisson coefficients

$r, \theta$  - spherical coordinates

$\mathbf{k}$  - flow direction unit vector







# Stress state

$$\varepsilon_{\alpha\beta} = \begin{cases} \varepsilon_r = \frac{\partial u_r}{\partial r} \\ \varepsilon_\theta = \frac{1}{r} \left( \frac{\partial u_\theta}{\partial \theta} + u_r \right) \\ \varepsilon_\varphi = \frac{1}{r} (u_\theta \operatorname{ctg} \theta + u_r) \\ \gamma_{r\theta} = \frac{1}{r} \frac{\partial u_r}{\partial \theta} + r \frac{\partial}{\partial r} \left( \frac{u_\theta}{r} \right) \end{cases}$$

$$\sigma_\alpha = 2G \left\{ \varepsilon_\alpha + \frac{\nu \sum_{i=1}^3 \varepsilon_i}{1-2\nu} \right\}$$

$$\tau_{\alpha\beta} = G\gamma_{\alpha\beta}; \quad \alpha, \beta = r, \theta, \varphi, \quad \alpha \neq \beta$$

# Stress intensity

$$\tau_i = \sqrt{I_2} = \sqrt{\frac{1}{6} \left[ \sigma_r - \sigma_\theta \right]^2 + \left[ \sigma_r - \sigma_\varphi \right]^2 + \left[ \sigma_\theta - \sigma_\varphi \right]^2 + \tau_{r\theta}^2}$$

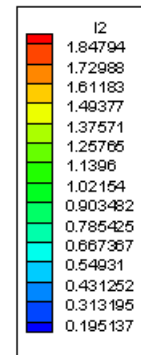
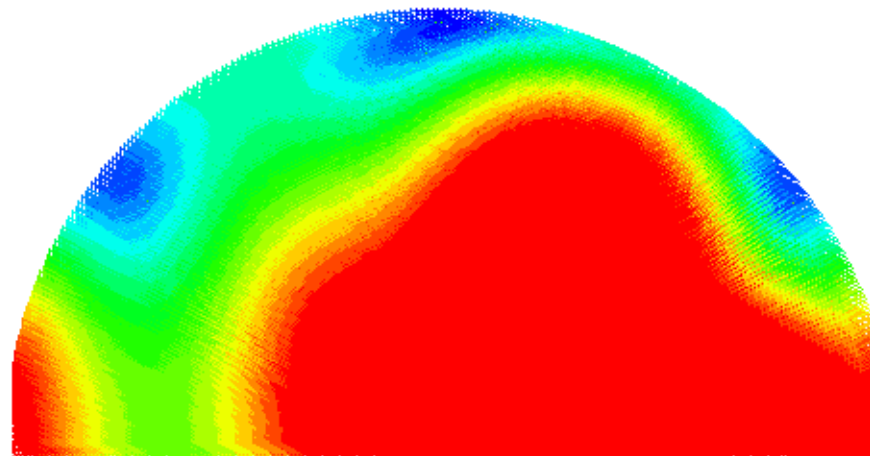
## Solution of problem:

L.A.Egorova The stress–strain state and disintegration of a meteoroid moving through the atmosphere/Journal of Applied Mathematics and Mechanics, 2011,

DOI: 10.1016/j.jappmathmech.2011.07.015



# Stress intensity





**The condition of safety: the dense of shape change potential energy less then critical one**

$$\sqrt{3}\tau_i = \sqrt{3}\sqrt{I_2} < \sigma_*$$

**If  $V \cong const$ , then  $I_2=I_2(\rho_2)$**

$$\frac{\rho_2(h)}{\rho_2(H)} = e^{-\frac{h}{H}}$$

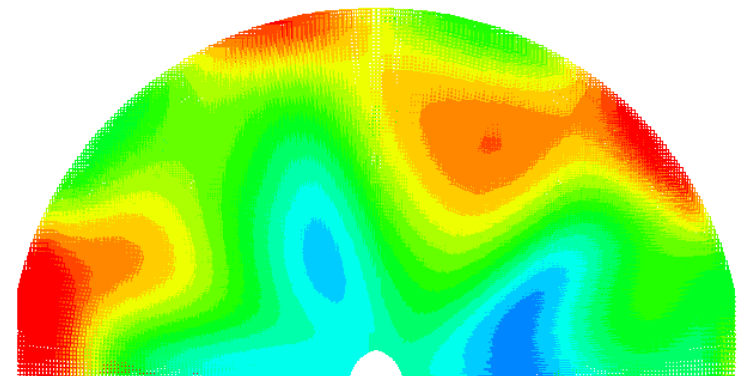
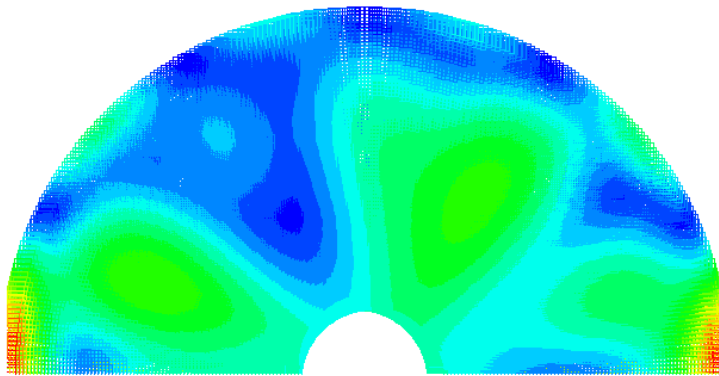
$$h = -H \ln \frac{\sigma_*}{\sqrt{3} \sqrt{I_2(H)}_{max}}$$

$\sigma_* = 700 \text{ atm}$ ,  $V = 30 \text{ km/s}$ ,  $h_{start} = 7,868 \text{ km}$  Falls to the surface

$\sigma_* = 50 \text{ atm}$ ,  $V = 30 \text{ km/s}$ ,  $h_{start} = 37,6 \text{ km}$ ,  $h_{end} = 26,6 \text{ km}$

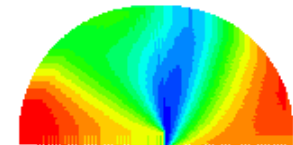
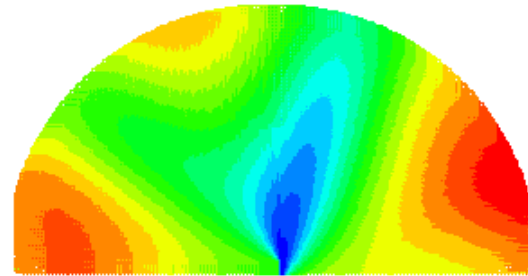
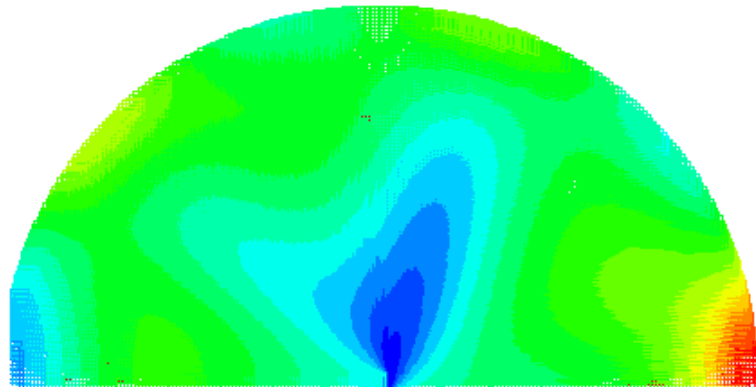


# Thermoelastic stress





# Problem of vanishing ball (Stephen problem)





# Criteria of elastic fragmentation

Condition of no rising flow strength

$$\frac{d}{dt} (\rho V^2) \geq 0, \quad \text{где} \quad \rho = \rho_0 \exp\left(-\frac{h - V_* t + at^2 / 2}{H}\right)$$

$$V(t) = V - at, \quad a = \frac{3V_*^2 \rho_0}{8r \rho_T} \exp\left(-\frac{h_*}{H}\right)$$

$$r \leq \frac{3}{4} \frac{H \rho_0}{\rho_T} e^{-\frac{h_*}{H}}$$

No fragmentation occur  
due to elastic stress

$$h_* = -H \ln \frac{\sigma_* 2\sqrt{2}}{\alpha \rho_0 V^2 \sqrt{\Sigma_{\max}}}$$



# Spectrum of pieces for the body crashed by explosion

$$\frac{dN_m}{dm} = Cm^{\frac{k}{3}-2}, k = 1.2$$

$N_m$  – number of particles of mass  $m$

$C$  – constant

Solution

$$N_m = \frac{2}{3} \left( \frac{1}{m^{0.6}} - 1 \right), \quad \bar{m} = \frac{m}{M}$$

Nemtchinov I.V., Popova O.P., Teterev A.V.  
Entering the large meteoroids into atmosphere: theory and practice. // Inzhen. Phis. Zhurn. 1999. V. 72, 6. Pp. 1233 - 1265. (in Russian)



# Thermal explosion effect

$$\begin{cases} m \frac{dV}{dt} = \frac{1}{2} c_x \rho_g V^2 S \\ i^* \frac{dm}{dt} = \frac{1}{2} c_H \rho_g V^3 S \end{cases}$$

The particle will light until velocity  $> V_*$

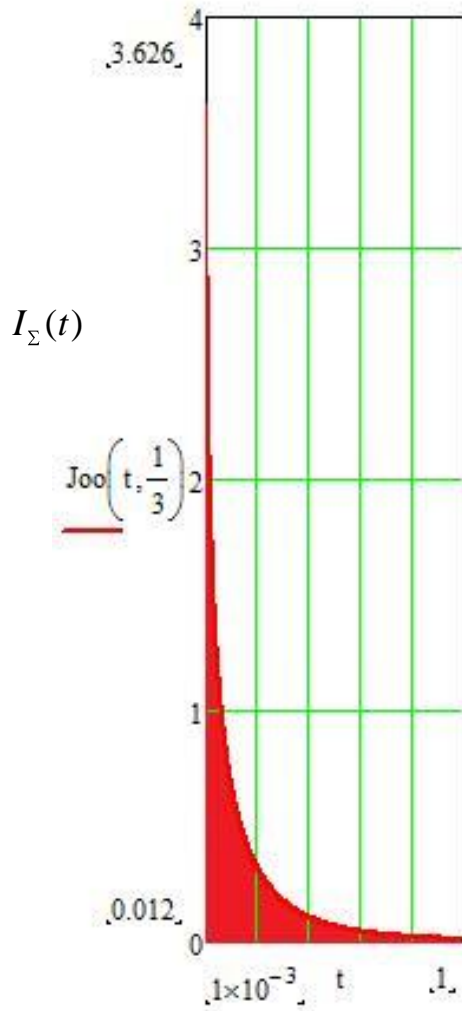
$$L = \int_0^{t_*} V dt = \frac{r_0}{A} \ln \left( 1 + \frac{V_0 - V_*}{V_*} \right), \quad A = \frac{3}{8} c_x \frac{\rho_g}{\rho_b}$$

$$t_* = \frac{8(V_0 - V_*)\rho_b}{3c_x \rho_g V_0 V_*} R$$





# Luminosity



$$I = -\tau \frac{V^2}{2} \frac{dM}{dt}$$

$$I_\Sigma(t) = \int_{m_*}^1 N_{m_0} \frac{d}{dm_0} \left( -\tau \frac{V^2}{2} \frac{dm}{dt} \right) dm_0$$



# Conclusions

**The body moving in the planet atmosphere is under the influence of the aerodynamic loads, the forces of inertia and the heat flux. As a result, the body undergoes ablation and even could be completely destroyed.**

**First of all, the stressed state within the body at any time is determined through an accurate solution of the Lamé equations.**

**During the flight of small fragments the thermo elastic forces become significant.**

**Finally, «thermal explosion» due to the rapid evaporation of small fragments cloud with a typical range of sizes of fragments was considered.**

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